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# Isotope exchange reactions involving $\text{HCO}^+$ with CO: A theoretical approach

M. Mladenović<sup>1</sup> and E. Roueff<sup>2</sup>

<sup>1</sup> Université Paris-Est, Laboratoire Modélisation et Simulation Multi Echelle, MSME UMR 8208 CNRS, 5 bd Descartes, 77454 Marne-la-Vallée, France

e-mail: Mirjana.Mladenovic@u-pem.fr

<sup>2</sup> LERMA, Observatoire de Paris, PSL Research University, CNRS, Sorbonne Universités, UPMC Univ. Paris 06, 92190 Meudon, France

e-mail: evelyne.roueff@obspm.fr

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## ABSTRACT

**Aims.** We aim to investigate fractionation reactions involved in the  $^{12}\text{C}/^{13}\text{C}$ ,  $^{16}\text{O}/^{18}\text{O}$ , and  $^{17}\text{O}$  balance.

**Methods.** Full-dimensional rovibrational calculations were used to compute numerically exact rovibrational energies and thermal equilibrium conditions to derive the reaction rate coefficients. A nonlinear least-squares method was employed to represent the rate coefficients by analytic functions.

**Results.** New exothermicities are derived for 30 isotopic exchange reactions of  $\text{HCO}^+$  with CO. For each of the reactions, we provide the analytic three-parameter Arrhenius-Kooij formula for both the forward reaction and backward reaction rate coefficients, that can further be used in astrochemical kinetic models. Rotational constants derived here for the  $^{17}\text{O}$  containing forms of  $\text{HCO}^+$  may assist detection of these cations in outer space.

**Key words.** ISM: general – ISM: molecules – ISM: abundances

## 1. Introduction

The  $^{18}\text{O}$  and  $^{17}\text{O}$  isotopic variants of CO are routinely detected in interstellar galactic and extragalactic environments and are used to determine the evolution trend of the corresponding  $^{18}\text{O}/^{17}\text{O}$  ratio through the galactic disk. However,  $\text{HC}^{17}\text{O}^+$  has, to our knowledge, only been detected in two sources, SgB2 by Guélin et al. (1982) and in the molecular peak of the L1544 pre-stellar core by Dore et al. (2001a), who also refined the spectroscopic constants and the hyperfine coupling constants as the  $^{17}\text{O}$  nucleus has a spin of  $5/2$ . With the advent of sensitive receivers and large collecting areas available in modern observational facilities, such as ALMA and NOEMA, there is no doubt that more observations of these molecular ions will become available. However, the chemical reactions which may occur and possibly enhance the abundance of this rare molecular ion through isotopic exchange reactions, such as those occurring for  $\text{HC}^{18}\text{O}^+$  (Mladenović & Roueff 2014), have not yet been reported. The purpose of the present study is to derive accurate values of the exothermicities involved in isotopic exchange reactions and to propose the corresponding reaction rate coefficients which can further be used in astrochemical models.

The fractionation of stable isotopes can be ascribed to a combination of the mass dependent thermodynamic (equilibrium) partition functions, the mass dependent diffusion coefficients, and the mass dependent reaction rate coefficients. This is in accordance with quantum mechanics, which predicts that mass affects the strength of chemical bonds and the vibrational, rotational, and translational motions, so that temperature dependent isotope fractionations may arise from quantum mechanical effects on rovibrational motion. For a given

vibrational state, the vibrational energy is lower for a bond involving a heavier isotope. The extent of isotopic fractionation varies inversely with temperature and is large at low temperatures.

Smith & Adams (1980) measured the forward  $k_f$  and reverse  $k_r$  rate coefficients for three isotopic variants of the reaction  $\text{HCO}^+$  with CO at 80, 200, 300, and 510 K using a selected ion flow tube (SIFT) technique. Langer et al. (1984) extrapolated the experimental values of Smith & Adams (1980) to temperatures below 80 K toward the limit of the average dipole orientation model of ion-polar molecule capture collisions (Su & Bowers 1975), producing the total rate coefficients  $k_T = k_f + k_r$  for nine temperatures over the range 5–300 K. Langer et al. employed a common reduced mass for three isotopic variants of  $\text{HCO}^+$ +CO studied by Smith & Adams (1980). The values for  $k_T$  were used in combination with theoretical spectroscopic parameters calculated for the isotopic variants of  $\text{HCO}^+$  by Henning et al. (1977) in order to model cosmochemical carbon and oxygen isotope fractionations. From the total mass-independent rate coefficients  $k_T$  and the theoretical zero point energy differences  $\Delta E$ , Langer et al. (1984) estimated the rate coefficients  $k_f$  and  $k_r$  for six reactions  $\text{HCO}^+$ +CO involving isotopologues containing  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ , and  $^{18}\text{O}$ .

Recently, we have investigated in some detail the isotope fractionation reactions of  $\text{HCO}^+/\text{HOC}^+$  with CO and of  $\text{N}_2\text{H}^+$  with  $\text{N}_2$  (Mladenović & Roueff 2014), hereafter called Paper I. In Paper I, we employed the global three-dimensional potential energy surfaces developed by Mladenović & Schmatz (1998) for the isomerizing system  $\text{HCO}^+/\text{HOC}^+$  and by Schmatz & Mladenović (1997) for the isoelectronic species  $\text{N}_2\text{H}^+$  in combination with a numerically exact method

for the rovibrational calculations (Mladenović & Bačić 1990; Mladenović 2002). For the reaction  $\text{HCO}^+ + \text{CO}$ , we pointed out inaccuracies of previous exothermicity values, which have been commonly used in chemical networks. The new exothermicities are found to affect significantly the rate coefficients derived at 10 K, corresponding to the temperature of dark interstellar cloud environments.

The isotopes H, D,  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ , and  $^{18}\text{O}$  were considered in our previous work (Mladenović & Roueff 2014), resulting in six reactions  $\text{HCO}^+ + \text{CO}$  involving hydrogen and six reactions  $\text{DCO}^+ + \text{CO}$  involving deuterium. In the present work, we extend our analysis with the stable isotope  $^{17}\text{O}$ . The possible isotopic variants of  $\text{HCO}^+$  and  $\text{CO}$  are connected by 15 reactions for the hydrogenic forms and 15 reactions for the deuterated counterparts. All 30 were studied here. Nominal abundances of oxygen isotopes  $^{16}\text{O}$ ,  $^{17}\text{O}$ , and  $^{18}\text{O}$  are 99.76, 0.04, and 0.20%, respectively (Mills et al. 1993).

Our theoretical approach is described in Sect. 2. In Sect. 3, we report the energies involved in all possible exchange reactions between  $\text{CO}$  and  $\text{HCO}^+$  ( $\text{HOC}^+$ ) isotopologues and their deuterium variants, as well as the rate coefficients for the reaction of  $\text{HCO}^+$  with  $\text{CO}$  in the 5–300 K temperature window. We summarize our finding in Sect. 4.

## 2. Methods

Isotopic exchange reactions can occur according to different types as discussed in Roueff et al. (2015). However, the equilibrium constant  $K_e$ , which provides the ratio of the forward reaction rate coefficient  $k_f$  and the backward (reverse) reaction rate coefficient  $k_r$ , is uniquely defined under thermal equilibrium conditions.

We considered the exchange between an heavy (H) and light (L) isotope in the reaction



Under thermal equilibrium conditions, the equilibrium constant  $K_e$  is given by

$$K_e = \frac{k_f}{k_r} = F_q e^{\Delta E/k_B T}, \quad (2)$$

using

$$F_q = f_m^{3/2} \frac{Q_{\text{int}}(\text{AH}) Q_{\text{int}}(\text{BL})}{Q_{\text{int}}(\text{AL}) Q_{\text{int}}(\text{BH})}, \quad (3)$$

and

$$f_m = \frac{m(\text{AH}) m(\text{BL})}{m(\text{AL}) m(\text{BH})}, \quad (4)$$

where  $m(\text{X})$  stands for the mass of the species X. In Eq. (2),  $\Delta E$  is the zero point energy difference between the reactants and the products,

$$\Delta E = E_0(\text{AL}) + E_0(\text{BH}) - E_0(\text{AH}) - E_0(\text{BL}). \quad (5)$$

To express  $\Delta E$  in Kelvin, we used  $\Delta E/k_B$ , where  $k_B$  is the Boltzmann constant. The zero point energy  $E_0(\text{X})$  for the species X is measured on an absolute energy scale. In practical applications,  $E_0(\text{X})$  is given relative to the respective potential energy minimum. The electronic states are not changed in the course of the reaction of Eq. (1), so that the ratio of the electronic partition functions is unity in Eq. (3). The term  $f_m$  arises from

the translational contribution. For the internal partition function  $Q_{\text{int}}$ , we used

$$Q_{\text{int}} = g \sum_J \sum_i (2J + 1) e^{-\varepsilon_i^J/k_B T}, \quad (6)$$

where  $\varepsilon_i^J = E_i^J - E_0^0$  denotes the rovibrational energy for a total angular momentum  $J$ . The factor  $2J + 1$  accounts for the degeneracy with respect to the space-fixed reference frame and  $g$  for the nuclear spin degeneracy. Additional care is required for the nuclear spin degeneracy factor when different spin states (e.g., ortho or para) are involved either in the reactants or in the products, as discussed in Terzieva & Herbst (2000). The rovibrational energies  $\varepsilon_i^J$  are measured relative to the corresponding zero point energy  $E_0^0$ .

For all the species involved in the reaction considered, the rovibrational energies  $\varepsilon_i^J$  are computed by theoretical means, considering explicitly the quantum mechanical effects due to vibrational anharmonicities and rovibrational couplings. The energies  $\varepsilon_i^J$  are used to evaluate the partition functions  $Q_{\text{int}}$  of Eq. (6) for a given temperature and then to compute the equilibrium constant  $K_e$  of Eq. (2). This approach has been pursued also in Paper I.

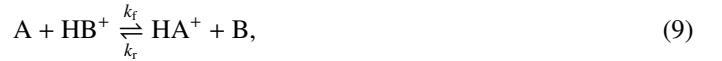
If the exchange proceeds through the formation of an adduct which can dissociate both backwards and forwards, we can define a total rate coefficient  $k_T$ , which is often expressed as the capture rate constant,

$$k_T = k_f + k_r. \quad (7)$$

We then readily have

$$k_f = k_T \frac{K_e}{K_e + 1} \quad \text{and} \quad k_r = k_T \frac{1}{K_e + 1}. \quad (8)$$

Such a mechanism takes place for the isotopic exchange reaction of  $^{13}\text{C}^+$  with  $\text{CO}$  and isotopologues as well as for proton transfer reactions of the type



where A and B are isotopologues of CO.

At low temperatures, the dominant term in the expression of Eq. (2) for the equilibrium constant  $K_e$  is the exponential term. Approximating  $F_q \approx 1$ , so that  $K_e \approx e^{\Delta E/k_B T}$ , it follows that

$$k_f \approx k_T \frac{1}{1 + e^{-\Delta E/k_B T}} \quad \text{and} \quad k_r \approx k_T \frac{e^{-\Delta E/k_B T}}{1 + e^{-\Delta E/k_B T}}. \quad (10)$$

The partition function factor  $F_q$  of Eq. (3) provides thus a quantitative estimate of the goodness of the approximation of Eq. (10), which is often employed in kinetic models at low temperatures.

## 3. Results and discussion

All the isotopic variants of  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ ,  $^{17}\text{O}$ , and  $^{18}\text{O}$  for  $\text{CO}$ ,  $\text{HCO}^+/\text{HOC}^+$ , and  $\text{DCO}^+/\text{DOC}^+$  are considered in this work. Using the global three-dimensional potential energy surface of Mladenović & Schmatz (1998) for the isomerizing system  $\text{HCO}^+/\text{HOC}^+$ , we calculate the rovibrational level energies for six isotopologues of  $\text{HCO}^+$  and six isotopologues of  $\text{DCO}^+$ , as well as for  $\text{HOC}^+$  and  $\text{DOC}^+$ . For isotopologues of  $\text{CO}$ , we employ the CCSD(T)/cc-pVQZ potential energy curve, developed previously (Mladenović & Roueff 2014). From these results, the zero point energy differences are readily determined for

the proton transfer reactions of Eq. (9) involving  $\text{HCO}^+$  with CO and  $\text{HOC}^+$  with CO. The equilibrium constants  $K_e$  as a function of temperature are evaluated according to Eq. (2) for the isotope exchange reactions  $\text{HCO}^+ + \text{CO}$ , which have already been studied experimentally (Smith & Adams 1980).

The spectroscopic properties for the isotopologues of  $\text{HCO}^+$  are summarized in Table A.1. The mode labels  $\nu_1, \nu_2, \nu_3$  refer respectively to the higher-frequency (CH) stretching vibration, the bending vibration, and the lower-frequency stretching (CO) vibration, whereby the bending vibration is doubly degenerate. By fitting the calculated ground state vibrational energies obtained for  $0 \leq J \leq 15$  to the standard polynomial expression,

$$E_0(J) = B_0 J(J+1) - D_0 [J(J+1)]^2, \quad (11)$$

we have derived the effective rotational constant  $B_0$  and the quartic centrifugal distortion constant  $D_0$  for the ground vibrational state. The ground state vibrational correction  $\Delta B_0$  to the equilibrium rotational constant  $B_e$  is given by

$$\Delta B_0 = B_e - B_0 = \frac{1}{2} (\alpha_1 + 2\alpha_2 + \alpha_3), \quad (12)$$

where  $\alpha_i$  is a vibration-rotation interaction constant for the  $i$ th vibration (Herzberg 1991). In this expression, we substitute our  $B_e$  values with the best estimate  $B_e^{\text{est}}$  of the equilibrium rotational constant, computed employing the best estimate of the equilibrium geometry  $r_e(\text{HC}) = 1.09197 \text{ \AA}$  and  $r_e(\text{CO}) = 1.10546 \text{ \AA}$  due to Puzzarini et al. (1996). This approach yields our best estimate  $B_0^{\text{est}}$  of the rotational constant for the ground vibrational state as  $B_0^{\text{est}} = B_e^{\text{est}} - \Delta B_0$ . As seen in Table A.1, all of the theoretical values for  $B_0^{\text{est}}$  agree with the available experimental  $B_0$  values within 11 MHz. Similar accuracy can also be expected for other  $\text{HCO}^+$  isotopologues that are not yet experimentally detected. For the fundamental vibrational transitions, our theoretical results agree within  $5 \text{ cm}^{-1}$  with the experimental findings. In the case of the quartic centrifugal distortion constants  $D_0$ , the agreement is within 4.5 kHz.

Direct proton transfer via the collinear approach from the carbon or oxygen side of CO to  $\text{HCO}^+/\text{HOC}^+$  was studied in Paper I using the CCSD(T)/aug-cc-pVTZ electronic structure method. In all the cases analyzed there, the reaction was found to proceed through a stable intermediate proton-bound complex  $\text{A-H-B}^+$  (see Eq. (9)), such that the reaction is barrierless in this description. In the case of the reaction involving  $\text{HOC}^+$ , it appears that the CO-catalyzed isomerisation is a more likely event than the proton transfer reaction, as seen in Fig. 2a of Paper I.

### 3.1. Zero point energies

The ground state vibrational energies calculated in this work for isotopic variants of CO,  $\text{HCO}^+$ , and  $\text{HOC}^+$  are collected in Table A.2. They are given relative to the corresponding potential energy minimum. On the potential energy surface for  $\text{HCO}^+/\text{HOC}^+$ , the global potential energy minimum is at  $-1.86 \text{ cm}^{-1}$ . In assembling Table A.2, we noticed that Table 1 of Paper I did not incorporate the energy shift of  $1.86 \text{ cm}^{-1}$  for the zero point energies of  $\text{D}^{12}\text{C}^{16}\text{O}^+$ ,  $\text{D}^{12}\text{C}^{18}\text{O}^+$ ,  $\text{D}^{13}\text{C}^{16}\text{O}^+$ , and  $\text{D}^{13}\text{C}^{18}\text{O}^+$ . Accordingly, we list the zero point energies for all of the isotopologues of CO,  $\text{HCO}^+$ , and  $\text{HOC}^+$  in Table A.2 of this work. We note that the other results of Paper I are not affected by this misprint.

The zero point energies are graphically displayed in Fig. 1 for the isotopic  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ ,  $^{17}\text{O}$ , and  $^{18}\text{O}$  variants of CO and  $\text{HCO}^+$ . The largest difference  $\Delta E_{\text{max}}$  seen there is between the zero point energies for the lightest and heaviest forms, yielding

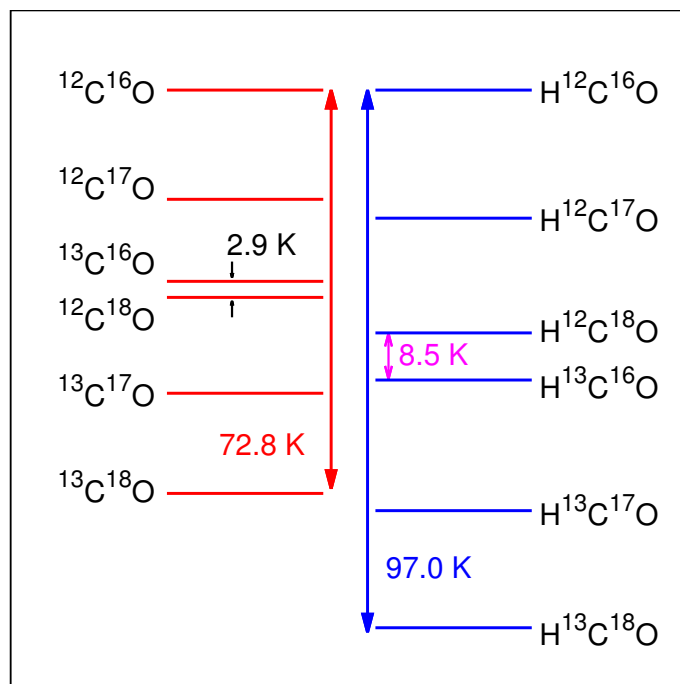


Fig. 1. Relative positions of the ground state vibrational energies of six isotopologues of CO and  $\text{HCO}^+$ .

$\Delta E_{\text{max}} = 72.8 \text{ K}$  for carbon monoxide and  $\Delta E_{\text{max}} = 97.0 \text{ K}$  for the formyl cation. The species containing  $^{17}\text{O}$  possess zero point energies lying between the zero point energies of the  $^{18}\text{O}$  and  $^{16}\text{O}$  forms in Fig. 1 and Table A.2.

### 3.2. Reactions of $\text{HCO}^+$ and $\text{HOC}^+$ with CO

The zero point energy differences for the proton transfer reactions  $\text{CO} + \text{HCO}^+/\text{HOC}^+$  are presented in Table A.3. In accordance with the notation of Paper I, the reactions involving the formyl cation  $\text{HCO}^+$  are labeled with F and the reactions involving the isoformyl cation  $\text{HOC}^+$  with I. The deuterium variant of the reaction Fn is denoted by Fn(D), where  $n = 1-15$ , and similar for the other cases. The reactions F1-F6 involving the isotopes  $^{16}\text{O}$  and  $^{18}\text{O}$  were studied in some detail in Paper I. Inclusion of the isotope  $^{17}\text{O}$  leads to nine additional reactions, which are denoted by F7-F15 and similar for the other variants. A complete list of possible reactions for the isotopes H, D,  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ ,  $^{17}\text{O}$ , and  $^{18}\text{O}$  is given in Table A.3. The exothermicities for the reactions F1-F6, F1(D)-F6(D), I1-I6, and I1(D)-I6(D) were already shown in Table 2 of Paper I. The proton transfer reactions I1 and I1(D) involving  $\text{HOC}^+$  were also considered by Lohr (1998).

The largest  $\Delta E$  in Table A.3 is associated with the reactions F5 and F5(D) in the case of  $\text{HCO}^+$  and with the reactions I6 and I6(D) in the case of  $\text{HOC}^+$ . From Table A.3, we easily deduce that  $^{13}\text{C}$  is preferentially placed in  $\text{H}/\text{DCO}^+$  and a heavier O in  $\text{H}/\text{DOC}^+$ . The reactions involving the same isotope of C in  $\text{HCO}^+$  and CO possess smaller exothermicities than the reactions involving different C isotopes, as seen by comparing for example, reaction F3 ( $\Delta E = 6.4 \text{ K}$ ) with reaction F5 ( $\Delta E = 24.2 \text{ K}$ ), which involve  $^{18}\text{O}$ . The corresponding  $^{17}\text{O}$  counterparts have somewhat smaller  $\Delta E$  values, for example reaction F7 of  $\Delta E = 3.4 \text{ K}$  versus reaction F11 of  $\Delta E = 21.2 \text{ K}$ . In all cases, reactions involving deuterium possess slightly higher exothermicities, for example,  $\Delta E$  for reactions F15 and F15(D)



are 20.8 K and 25.3 K, respectively. The reactions with the isoformyl isomers have lower exothermicities than the reactions with the formyl forms. In several cases, the reactions  $\text{HOC}^+ + \text{CO}$  proceed in the direction opposite to the direction of the corresponding  $\text{HCO}^+ + \text{CO}$  reaction in accordance with the fact that the isotopic substitution of the central atom is thermodynamically favored. Table A.3 clearly exemplifies this effect.

The partition function factors  $F_q$ , the isotope exchange equilibrium constants  $K_e$ , and the rate coefficients ( $k_f, k_r$ ) for several temperatures between 5 and 300 K are given in Table A.4 for reactions F7-F15 and F7(D)-F15(D). This table complements Table 5 of Paper I, which provides analogous information for reactions F1-F6 and F1(D)-F6(D). The values of  $K_e$  are obtained using Eq. (2) by direct evaluation of the internal partition functions  $Q_{\text{int}}$  from the computed rovibrational energies. The forward reaction  $k_f$  and backward reaction  $k_r$  rate coefficients are calculated according to Eq. (8) using our  $\Delta E$  values of Table A.2 in combination with the total temperature dependent rate coefficients  $k_T$  given by Langer et al. (1984).

In Eq. (2),  $F_q$  is a mass and temperature dependent factor, defined by Eq. (3). Its temperature dependence is due to the temperature dependence of  $Q_{\text{int}}$ . The mass dependence of  $F_q$  comes from the translational contribution  $f_m$ . In addition, the mass affects the effective rotational constants for a given vibrational state, as well as the reduced mass specifying the vibrational motion, yielding thus the mass dependent  $Q_{\text{int}}$ . In the low temperature limit relevant for dark cloud conditions, the discrete rotational structure of the ground vibrational state provides the major contribution to  $Q_{\text{int}}$ . For 30 reactions  $\text{HCO}^+ + \text{CO}$ , given here in Table A.4 and before in Table 5 of Paper I, the factor  $F_q$  differs from 1 at most by 3.5%. Using  $F_q = 1$  to compute ( $k_f, k_r$ ) by means of Eq. (10), we obtain rate coefficients which differ by at most 4% from the values listed in Table A.4.

### 3.3. Analytic representations of the rate coefficients

For the reaction  $\text{HCO}^+ + \text{CO}$ , the total rate coefficients derived by Langer et al. (1984) are available for nine temperatures: 5, 10, 20, 40, 60, 80, 100, 200, and 300 K. Using these values, the forward and backward rate coefficients are determined (see Table A.4) and fit via the popular Arrhenius-Kooij formula (Kooij 1893),

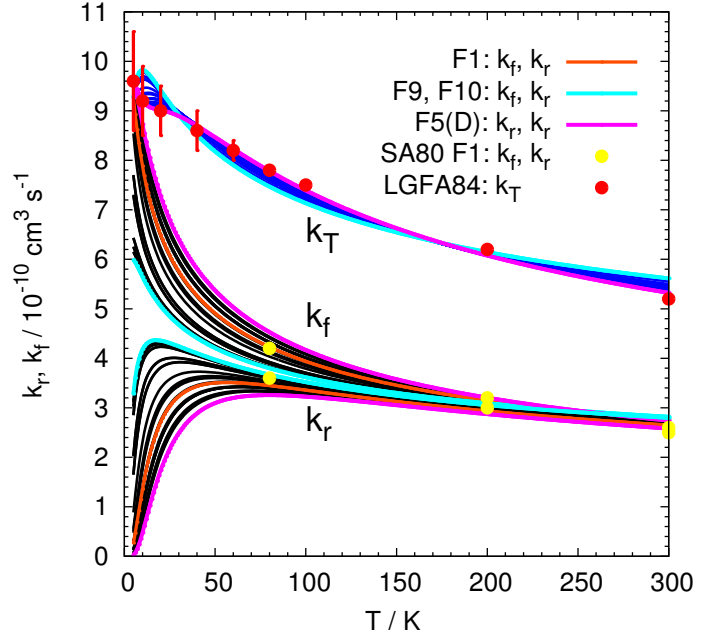
$$k(T) = A \left( \frac{T}{T_{\text{ref}}} \right)^b e^{-C/T}, \quad \text{where} \quad T_{\text{ref}} = 298 \text{ K}. \quad (13)$$

This analytical expansion has one nonlinear parameter, so that we employ a nonlinear least-squares technique (the Levenberg-Marquardt algorithm) to obtain optimum values for the fitting parameters (Press et al. 1985). The resulting values for  $A$ ,  $b$ , and  $C$  are given in Table A.5. The statistical uncertainties are given as root-mean-square (rms) fitting errors,

$$\text{rms} = \sqrt{\frac{\sum_{i=1}^N (y_i - f_i)^2}{N}}, \quad (14)$$

where  $N$  is the number of the known (input) data ( $x_i, y_i$ ) fit by a function  $f = f(x)$ , so that  $f_i = f(x_i)$ .

The forward reaction rate coefficients  $k_f$  and the backward reaction rate coefficients  $k_r$  are fit separately since they represent elementary chemical processes. The difference between the parameters  $C_r$  for  $k_r$  and  $C_f$  for  $k_f$  is equal to the exothermicity for the corresponding isotopic exchange reaction,  $C_r - C_f = \Delta E$ .



**Fig. 2.** Temperature dependence of the rate coefficients  $k_T$ ,  $k_f$ , and  $k_r$  obtained in the fitting procedures.

This is easy to verify by comparing Tables A.3 and A.5. The fitting parameters in Table A.5 reproduce the values of the rate coefficients  $k_f$  and  $k_r$  at nine temperatures (Table A.4 of this work and Table 5 of Paper I) with rms deviations better than  $2 \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ . Mean absolute percentage deviations are better than 4%.

The variation with temperature of the rate coefficients  $k_f$  and  $k_r$  obtained in the fitting procedures are graphically displayed in Fig. 2 for all 30 isotopic variants of the reaction  $\text{HCO}^+ + \text{CO}$ . At temperatures above 50 K,  $k_f$  and  $k_r$  exhibit a weak temperature dependence. The forward reaction becomes faster and the backward reaction slower as the temperature decreases, so that  $k_r$  tends to zero as  $T \rightarrow 0 \text{ K}$ . In Fig. 2,  $k_f$  and  $k_r$  are most different for reaction F5(D), associated with the largest  $\Delta E$  value of 29.4 K in Table A.3. They are least different for reactions F9 and F10, attributed the smallest  $\Delta E$  value of 3.0 K (Table A.3). The experimental results of Smith & Adams (1980) available for the reaction F1 at 80, 200, and 300 K are also shown (yellow circles).

In Fig. 2, the total rate coefficients  $k_T$  (blue lines) are obtained as  $k_f + k_r$  for each of the 30 reactions considered. The values of Langer et al. (1984) (red circles) are additionally shown along with their estimated uncertainties (vertical bars). Even though  $k_f$  and  $k_r$  are noticeably different for different reactions, the resulting  $k_T$  functions are similar, as expected from the model used. Our  $k_T$  results for reaction F5(D) best approximate the values of Langer et al. (1984) and can be used as an analytic representation for their values if/when needed.

The estimates of Langer et al. (1984) cover temperatures between 5 and 300 K. The analytic expressions of Table A.5 are accordingly valid only over this temperature range. Since the modified Arrhenius function of Eq. (13) goes always to zero when  $T \rightarrow 0 \text{ K}$ , the functional representations derived for  $k_f$  will also tend to zero at temperatures below 5 K. Prior to elaborating on other forms more suitable for kinetic applications close to 0 K, one may consider the inclusion in astrochemical kinetic networks of the analytic three-parameter representations for the

rate coefficients  $k_f$  and  $k_r$ , derived here for the isotope exchange reactions  $\text{HCO}^+ + \text{CO}$  (Table A.5).

#### 4. Summary

In the present work, we have employed the full-dimensional quantum mechanical methods to calculate the rovibrational energies for all isotopologues of CO,  $\text{HCO}^+$ , and  $\text{HOC}^+$  involving the isotopes H, D,  $^{16}\text{O}$ ,  $^{17}\text{O}$ ,  $^{18}\text{O}$ ,  $^{12}\text{C}$ , and  $^{13}\text{C}$ . These results are used to derive accurate values of the exothermicities for possible isotopic exchange reactions. For the reaction of  $\text{HCO}^+$  with CO, all possible isotope fractionation variants are subsequently considered (in total 30 reactions). Values corresponding to the  $^{17}\text{O}$  isotope are reported for the first time. The energy defects arising for the  $^{17}\text{O}$  cases are found to be slightly smaller than those involved with  $^{18}\text{O}$ .

For each of the reactions considered, the analytic three-parameter expressions are derived for the isotopic exchange rate coefficients  $k_f$  and  $k_r$ . These analytic representations can straightforwardly be introduced in astrochemical kinetic models in order to better understand the isotopic chemical evolution.

For all of the isotopologues of  $\text{HCO}^+$  involving H, D,  $^{16}\text{O}$ ,  $^{17}\text{O}$ ,  $^{18}\text{O}$ ,  $^{12}\text{C}$ , and  $^{13}\text{C}$ , we provide the fundamental vibrational transitions and the rotational constants. Our best estimates of the rotational constants can provide useful assistance in analyzing expected observations of the rare forms of this cation.

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## Appendix A: Tables

**Table A.1.** Computed anharmonic fundamental vibrational transitions  $\nu_i$  (in  $\text{cm}^{-1}$ ), estimated ground state vibrational corrections  $\Delta B_0$  (in MHz), best estimates of the rotational constant  $B_0^{\text{est}}$  (in MHz), and quartic centrifugal distortion rotational constants  $D_0$  (in kHz) for the isotopologues of  $\text{HCO}^+$ .

Species	$\nu_1$	$\nu_2$	$\nu_3$	$\Delta B_0$	$B_0^{\text{est}}$	$D_0$
$\text{H}^{12}\text{C}^{16}\text{O}^+$	3085.6 [3088.7] <sup>a</sup>	830.7 [829.7] <sup>b</sup>	2179.1 [2183.9] <sup>c</sup>	243.7	44 605.1 [44 594.4] <sup>d</sup>	81.9 [82.4] <sup>e</sup>
$\text{H}^{12}\text{C}^{17}\text{O}^+$	3083.1	829.6	2152.3	235.4	43 539.2 [43 528.9] <sup>f</sup>	78.1 [79.1] <sup>f</sup>
$\text{H}^{12}\text{C}^{18}\text{O}^+$	3081.0	828.7	2128.2	228.1	42 591.2 [42 581.3] <sup>d</sup>	74.7 [76] <sup>d</sup>
$\text{H}^{13}\text{C}^{16}\text{O}^+$	3063.0	823.2	2145.4	234.1	43 387.6 [43 377.3] <sup>d</sup>	77.6 [79] <sup>d</sup>
$\text{H}^{13}\text{C}^{17}\text{O}^+$	3060.9	822.1	2117.5	225.8	42 305.6	73.8
$\text{H}^{13}\text{C}^{18}\text{O}^+$	3059.1	821.2	2092.5	218.5	41 343.1 [41 333.6] <sup>d</sup>	70.5 [75] <sup>d</sup>
$\text{D}^{12}\text{C}^{16}\text{O}^+$	2580.5 [2584.6] <sup>d</sup>	667.5	1900.8 [1904.1] <sup>d</sup>	173.5	36 027.6 [36 019.8] <sup>d</sup>	55.2 [55.8] <sup>e</sup>
$\text{D}^{12}\text{C}^{17}\text{O}^+$	2566.3	666.1	1885.3	167.6	35 178.4	52.5
$\text{D}^{12}\text{C}^{18}\text{O}^+$	2554.2	665.0	2554.2	162.4	34 421.1 [34 413.7] <sup>g</sup>	50.1
$\text{D}^{13}\text{C}^{16}\text{O}^+$	2529.5	658.0	1893.8	168.8	35 374.4 [35 366.6] <sup>g</sup>	52.8
$\text{D}^{13}\text{C}^{17}\text{O}^+$	2515.8	656.7	1877.2	162.8	34 508.7	50.1
$\text{D}^{13}\text{C}^{18}\text{O}^+$	2504.1	655.4	1861.7	157.5	33 736.4	47.8

**Notes.** Experimental values are given in brackets.

**References.** <sup>(a)</sup> Amano (1983). <sup>(b)</sup> Davies & Rothwell (1984). <sup>(c)</sup> Foster et al. (1984). <sup>(d)</sup> Taken from Puzzarini et al. (1996). <sup>(e)</sup> Hirota & Endo (1988). <sup>(f)</sup> Dore et al. (2001b). <sup>(g)</sup> As given at <http://physics.nist.gov/cgi-bin/MolSpec/triperiodic.pl>

**Table A.2.** Zero point vibrational energies (in  $\text{cm}^{-1}$ ) of the isotopologues of CO,  $\text{HCO}^+$ , and  $\text{HOC}^+$ .

Species	$E_0$	Species	$E_0$
$^{12}\text{C}^{16}\text{O}$	1079.11		
$^{12}\text{C}^{17}\text{O}$	1065.41		
$^{12}\text{C}^{18}\text{O}$	1053.11		
$^{13}\text{C}^{16}\text{O}$	1055.12		
$^{13}\text{C}^{17}\text{O}$	1041.09		
$^{13}\text{C}^{18}\text{O}$	1028.52		
$\text{H}^{12}\text{C}^{16}\text{O}^+$	3524.60	$\text{H}^{16}\text{O}^{12}\text{C}^+$	2871.08
$\text{H}^{12}\text{C}^{17}\text{O}^+$	3508.53	$\text{H}^{17}\text{O}^{12}\text{C}^+$	2855.41
$\text{H}^{12}\text{C}^{18}\text{O}^+$	3494.15	$\text{H}^{18}\text{O}^{12}\text{C}^+$	2841.37
$\text{H}^{13}\text{C}^{16}\text{O}^+$	3488.24	$\text{H}^{16}\text{O}^{13}\text{C}^+$	2848.66
$\text{H}^{13}\text{C}^{17}\text{O}^+$	3471.85	$\text{H}^{17}\text{O}^{13}\text{C}^+$	2832.71
$\text{H}^{13}\text{C}^{18}\text{O}^+$	3457.16	$\text{H}^{18}\text{O}^{13}\text{C}^+$	2818.42
$\text{D}^{12}\text{C}^{16}\text{O}^+$	2946.08	$\text{D}^{16}\text{O}^{12}\text{C}^+$	2357.61
$\text{D}^{12}\text{C}^{17}\text{O}^+$	2929.54	$\text{D}^{17}\text{O}^{12}\text{C}^+$	2340.97
$\text{D}^{12}\text{C}^{18}\text{O}^+$	2914.72	$\text{D}^{18}\text{O}^{12}\text{C}^+$	2326.06
$\text{D}^{13}\text{C}^{16}\text{O}^+$	2907.02	$\text{D}^{16}\text{O}^{13}\text{C}^+$	2334.87
$\text{D}^{13}\text{C}^{17}\text{O}^+$	2890.15	$\text{D}^{17}\text{O}^{13}\text{C}^+$	2317.97
$\text{D}^{13}\text{C}^{18}\text{O}^+$	2875.03	$\text{D}^{18}\text{O}^{13}\text{C}^+$	2302.82





**Table A.4.** Equilibrium constants  $K_e$ , partition function factors  $F_q$ , and rate coefficients  $k_f, k_r$  (in  $10^{-10} \text{ cm}^3 \text{ s}^{-1}$ ) for the reactions of H/DCO<sup>+</sup> with CO involving the isotope <sup>17</sup>O.

Reaction	T/K	5	10	20	40	60	80	100	200	300
F7	$F_q$	0.9981	0.9974	0.9971	0.9970	0.9969	0.9969	0.9969	0.9969	0.9970
	$K_e$	1.97	1.40	1.18	1.09	1.05	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.36, 3.24)	(5.37, 3.83)	(4.87, 4.13)	(4.48, 4.12)	(4.21, 3.99)	(3.98, 3.82)	(3.81, 3.69)	(3.12, 3.08)	(2.61, 2.59)
F8	$F_q$	0.9981	0.9974	0.9971	0.9969	0.9969	0.9969	0.9969	0.9969	0.9970
	$K_e$	1.97	1.40	1.18	1.09	1.06	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.37, 3.23)	(5.37, 3.83)	(4.88, 4.12)	(4.48, 4.12)	(4.21, 3.99)	(3.98, 3.82)	(3.81, 3.69)	(3.12, 3.08)	(2.61, 2.59)
F9	$F_q$	0.9983	0.9977	0.9974	0.9973	0.9972	0.9972	0.9972	0.9972	0.9973
	$K_e$	1.82	1.35	1.16	1.08	1.05	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.20, 3.40)	(5.28, 3.92)	(4.83, 4.17)	(4.46, 4.14)	(4.20, 4.00)	(3.97, 3.83)	(3.80, 3.70)	(3.12, 3.08)	(2.61, 2.59)
F10	$F_q$	0.9982	0.9977	0.9974	0.9973	0.9972	0.9972	0.9972	0.9972	0.9973
	$K_e$	1.83	1.35	1.16	1.08	1.05	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.21, 3.39)	(5.29, 3.91)	(4.84, 4.16)	(4.46, 4.14)	(4.20, 4.00)	(3.97, 3.83)	(3.80, 3.70)	(3.12, 3.08)	(2.61, 2.59)
F11	$F_q$	0.9835	0.9806	0.9793	0.9786	0.9784	0.9783	0.9782	0.9784	0.9795
	$K_e$	68.17	8.16	2.83	1.66	1.39	1.27	1.21	1.09	1.05
	$(k_f, k_r)$	(9.46, 0.14)	(8.20, 1.00)	(6.65, 2.35)	(5.37, 3.23)	(4.77, 3.43)	(4.37, 3.43)	(4.10, 3.40)	(3.23, 2.97)	(2.66, 2.54)
F12	$F_q$	0.9854	0.9832	0.9821	0.9816	0.9814	0.9813	0.9813	0.9815	0.9824
	$K_e$	34.67	5.83	2.39	1.53	1.32	1.23	1.17	1.07	1.04
	$(k_f, k_r)$	(9.33, 0.27)	(7.85, 1.35)	(6.35, 2.65)	(5.20, 3.40)	(4.67, 3.53)	(4.30, 3.50)	(4.05, 3.45)	(3.21, 2.99)	(2.65, 2.55)
F13	$F_q$	0.9870	0.9854	0.9847	0.9843	0.9841	0.9841	0.9840	0.9842	0.9851
	$K_e$	19.03	4.33	2.06	1.42	1.26	1.18	1.14	1.06	1.03
	$(k_f, k_r)$	(9.12, 0.48)	(7.47, 1.73)	(6.06, 2.94)	(5.05, 3.55)	(4.57, 3.63)	(4.23, 3.57)	(4.00, 3.50)	(3.19, 3.01)	(2.64, 2.56)
F14	$F_q$	0.9873	0.9857	0.9850	0.9846	0.9845	0.9844	0.9844	0.9845	0.9854
	$K_e$	17.57	4.16	2.02	1.41	1.25	1.18	1.14	1.06	1.03
	$(k_f, k_r)$	(9.08, 0.52)	(7.42, 1.78)	(6.02, 2.98)	(5.03, 3.57)	(4.56, 3.64)	(4.22, 3.58)	(3.99, 3.51)	(3.19, 3.01)	(2.64, 2.56)
F15	$F_q$	0.9836	0.9809	0.9795	0.9789	0.9787	0.9786	0.9785	0.9787	0.9798
	$K_e$	63.58	7.89	2.78	1.65	1.39	1.27	1.21	1.09	1.05
	$(k_f, k_r)$	(9.45, 0.15)	(8.16, 1.04)	(6.62, 2.38)	(5.35, 3.25)	(4.76, 3.44)	(4.36, 3.44)	(4.10, 3.40)	(3.23, 2.97)	(2.66, 2.54)
F7(D)	$F_q$	0.9969	0.9958	0.9953	0.9950	0.9950	0.9949	0.9949	0.9950	0.9953
	$K_e$	2.25	1.50	1.22	1.10	1.06	1.05	1.04	1.02	1.01
	$(k_f, k_r)$	(6.65, 2.95)	(5.52, 3.68)	(4.95, 4.05)	(4.51, 4.09)	(4.23, 3.97)	(3.99, 3.81)	(3.82, 3.68)	(3.12, 3.08)	(2.61, 2.59)
F8(D)	$F_q$	0.9968	0.9957	0.9952	0.9950	0.9949	0.9949	0.9948	0.9950	0.9953
	$K_e$	2.26	1.50	1.22	1.10	1.07	1.05	1.04	1.02	1.01
	$(k_f, k_r)$	(6.65, 2.95)	(5.52, 3.68)	(4.95, 4.05)	(4.51, 4.09)	(4.23, 3.97)	(3.99, 3.81)	(3.82, 3.68)	(3.12, 3.08)	(2.61, 2.59)
F9(D)	$F_q$	0.9971	0.9962	0.9957	0.9955	0.9955	0.9954	0.9954	0.9955	0.9958
	$K_e$	2.06	1.43	1.19	1.09	1.06	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.46, 3.14)	(5.42, 3.78)	(4.90, 4.10)	(4.49, 4.11)	(4.21, 3.99)	(3.98, 3.82)	(3.81, 3.69)	(3.12, 3.08)	(2.61, 2.59)
F10(D)	$F_q$	0.9970	0.9961	0.9957	0.9955	0.9954	0.9954	0.9954	0.9955	0.9958
	$K_e$	2.07	1.44	1.20	1.09	1.06	1.04	1.03	1.01	1.01
	$(k_f, k_r)$	(6.48, 3.12)	(5.42, 3.78)	(4.90, 4.10)	(4.49, 4.11)	(4.22, 3.98)	(3.98, 3.82)	(3.81, 3.69)	(3.12, 3.08)	(2.61, 2.59)
F11(D)	$F_q$	0.9733	0.9691	0.9671	0.9661	0.9658	0.9657	0.9656	0.9667	0.9692
	$K_e$	168.50	12.75	3.51	1.84	1.48	1.33	1.25	1.10	1.06
	$(k_f, k_r)$	(9.54, 0.06)	(8.53, 0.67)	(7.00, 2.00)	(5.57, 3.03)	(4.90, 3.30)	(4.46, 3.34)	(4.17, 3.33)	(3.25, 2.95)	(2.67, 2.53)
F12(D)	$F_q$	0.9763	0.9732	0.9717	0.9710	0.9707	0.9706	0.9706	0.9715	0.9738
	$K_e$	74.76	8.52	2.87	1.67	1.39	1.27	1.21	1.08	1.05
	$(k_f, k_r)$	(9.47, 0.13)	(8.23, 0.97)	(6.68, 2.32)	(5.38, 3.22)	(4.77, 3.43)	(4.37, 3.43)	(4.10, 3.40)	(3.22, 2.98)	(2.66, 2.54)
F13(D)	$F_q$	0.9791	0.9769	0.9759	0.9753	0.9752	0.9751	0.9750	0.9759	0.9779
	$K_e$	36.27	5.95	2.41	1.53	1.32	1.22	1.17	1.07	1.04
	$(k_f, k_r)$	(9.34, 0.26)	(7.88, 1.32)	(6.36, 2.64)	(5.20, 3.40)	(4.66, 3.54)	(4.29, 3.51)	(4.04, 3.46)	(3.20, 3.00)	(2.65, 2.55)
F14(D)	$F_q$	0.9795	0.9774	0.9764	0.9759	0.9757	0.9756	0.9756	0.9764	0.9784
	$K_e$	33.10	5.68	2.35	1.52	1.31	1.22	1.16	1.07	1.04
	$(k_f, k_r)$	(9.32, 0.28)	(7.82, 1.38)	(6.32, 2.68)	(5.18, 3.42)	(4.65, 3.55)	(4.28, 3.52)	(4.03, 3.47)	(3.20, 3.00)	(2.65, 2.55)
F15(D)	$F_q$	0.9734	0.9694	0.9675	0.9666	0.9663	0.9661	0.9661	0.9671	0.9697
	$K_e$	154.91	12.23	3.44	1.82	1.47	1.33	1.24	1.10	1.06
	$(k_f, k_r)$	(9.54, 0.06)	(8.50, 0.70)	(6.97, 2.03)	(5.55, 3.05)	(4.89, 3.31)	(4.45, 3.35)	(4.16, 3.34)	(3.24, 2.96)	(2.67, 2.53)

**Table A.5.** Fitting parameters  $A$  (in  $10^{-10} \text{ cm}^3 \text{ s}^{-1}$ ),  $b$ , and  $C$  (in K) for the forward rate coefficients  $k_f$  and the reverse rate coefficients  $k_r$  for the reactions of  $\text{HCO}^+$  with CO.

Reaction	$k_f$				$k_r$			
	$A_f$	$b_f$	$C_f$	rms	$A_r$	$b_r$	$C_r$	rms
F1	2.77	-0.33	0.64	0.08	2.82	-0.33	18.46	0.06
F2	2.77	-0.33	0.65	0.08	2.82	-0.33	18.51	0.06
F3	2.82	-0.25	0.32	0.19	2.83	-0.25	6.73	0.13
F4	2.82	-0.25	0.32	0.19	2.83	-0.25	6.78	0.13
F5	2.77	-0.37	1.31	0.06	2.84	-0.37	25.59	0.04
F6	2.79	-0.28	0.16	0.14	2.83	-0.28	11.56	0.09
F7	2.83	-0.23	0.91	0.20	2.84	-0.23	4.31	0.16
F8	2.83	-0.23	0.91	0.20	2.84	-0.23	4.32	0.16
F9	2.83	-0.23	1.02	0.20	2.84	-0.23	4.04	0.17
F10	2.83	-0.23	1.01	0.20	2.84	-0.23	4.06	0.17
F11	2.77	-0.35	0.99	0.07	2.83	-0.35	22.22	0.05
F12	2.77	-0.33	0.65	0.08	2.82	-0.33	18.48	0.06
F13	2.78	-0.31	0.37	0.10	2.82	-0.31	15.19	0.07
F14	2.78	-0.30	0.34	0.11	2.82	-0.30	14.75	0.07
F15	2.77	-0.35	0.96	0.07	2.83	-0.35	21.84	0.05
F1(D)	2.76	-0.35	1.03	0.06	2.84	-0.35	22.76	0.05
F2(D)	2.76	-0.35	1.03	0.06	2.84	-0.35	22.80	0.05
F3(D)	2.81	-0.26	0.19	0.18	2.83	-0.26	7.93	0.12
F4(D)	2.81	-0.26	0.19	0.17	2.83	-0.26	7.96	0.12
F5(D)	2.77	-0.39	1.80	0.04	2.87	-0.39	31.30	0.03
F6(D)	2.78	-0.30	0.29	0.11	2.83	-0.30	14.29	0.08
F7(D)	2.83	-0.24	0.73	0.20	2.84	-0.24	4.82	0.16
F8(D)	2.83	-0.24	0.73	0.20	2.84	-0.24	4.83	0.16
F9(D)	2.83	-0.24	0.84	0.20	2.84	-0.24	4.49	0.16
F10(D)	2.83	-0.24	0.84	0.20	2.84	-0.24	4.51	0.16
F11(D)	2.76	-0.37	1.45	0.05	2.85	-0.38	27.28	0.04
F12(D)	2.76	-0.35	1.03	0.06	2.84	-0.35	22.77	0.05
F13(D)	2.76	-0.33	0.66	0.08	2.83	-0.33	18.76	0.06
F14(D)	2.77	-0.33	0.61	0.08	2.83	-0.33	18.25	0.06
F15(D)	2.76	-0.37	1.41	0.05	2.85	-0.37	26.82	0.04

**Notes.** The rms fitting errors (in  $10^{-10} \text{ cm}^3 \text{ s}^{-1}$ ) are calculated using Eq. (14).